



Probing highly compressed degenerate matter and matter at extreme Gbar pressures at NIF

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NIF User Group
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Liaison scientist for NIF proposal, Postdoctoral Fellow

P. Neumayer, R. Falcone, D. Swift, J. Hawreliak, H. Lee, R. Redmer, E. Foerster, C. Fortmann, S. Le Pape, G. Hays, S. Rose, R. Hemley, R. Jeanloz, D. Hicks, P. Cellier, J. Eggert, D. Milathianaki, T. Doeppner, O. Landen, G. Collins, S. Glenzer

Matter at $>20\times$ solid compression and pressures of Gbar can only be reached and probed with XRTS and radiography at NIF

- **Matter at extreme densities occurs in astrophysical objects such as giant gas planets and highly evolved stars**
 - At extreme densities matter becomes metallic. The ions are strongly coupled due to the small inter-particle distance and the high charge state.
 - Electron coupling decreases at higher density, due to the increasing Fermi-energy
- **Matter at extreme pressures of Gbars (and higher temp) occurs in the cores of super-giant planets and stars**
 - Fundamental physics including EOS and ionization of condensed matter up to Gbar pressures is important for understanding the evolution of these astrophysical bodies

**Measurement of
microscopic properties
& plasma parameters**



**X-ray scattering &
X-ray radiography**

These joint proposals include many outside collaborators from several institutions

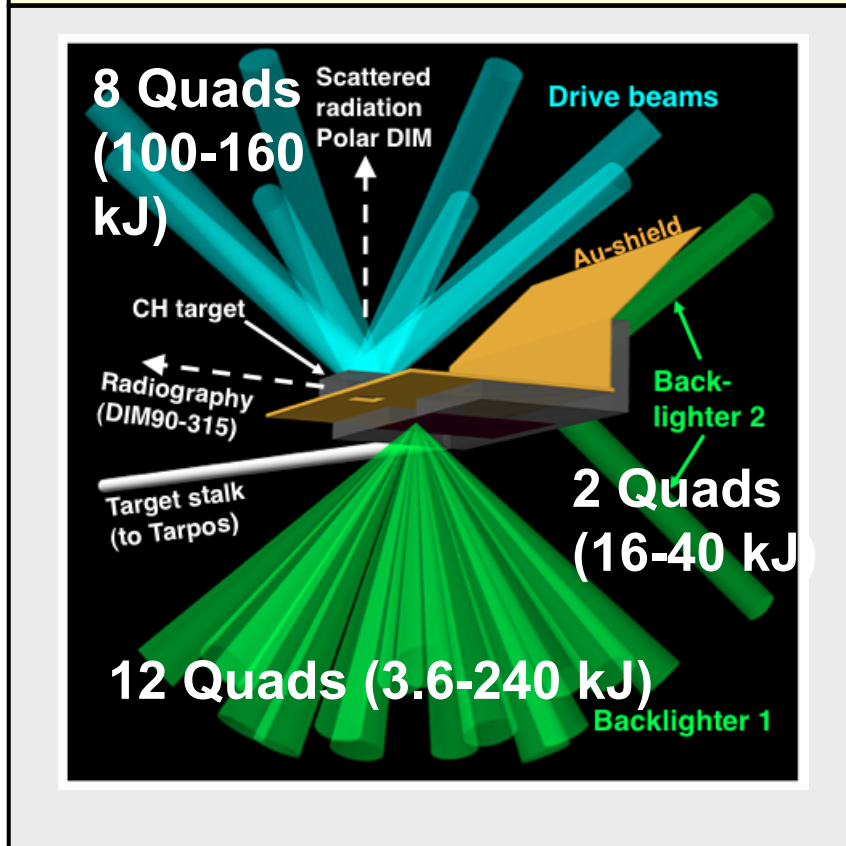


- GSI, Germany:
 - *P. Neumayer*
- Univ. of California Berkeley, USA/LBNL, USA:
 - *R. Falcone, R. Jeanloz*
- LLNL, USA
 - *D. Swift, J. Hawreliak, S. Le Pape, D. Hicks, P. Cellier, J. Eggert, T. Doeppner, O. Landen, G. Collins, S. Glenzer*
- LCLS, USA
 - *H. J. Lee, D. Milathianaki, G. Hays*
- Univ. of Jena, Germany
 - *E. Foerster, et al.*
- Univ. of Rostock, Germany
 - *R. Redmer, et al.*
- Univ. of California Los Angeles, USA
 - *C. Fortmann*
- Imperial College London
 - *S. Rose, et al.*
- Carnegie Institute of Washington
 - *R. Hemley, et al.*

Experiments at NIF will compress matter to densities of $>20\times$ solid, while staying on a low isentrope, using multiple shocks



CH will be directly driven in a planar geometry and probed with x-rays



A new spectrometer snout for XRTS is being developed...

(PI) P. Neumayer, et al.

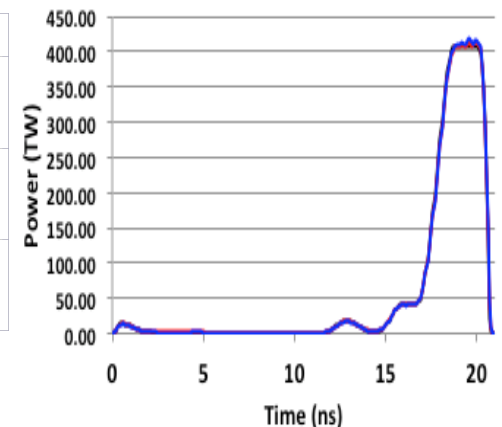
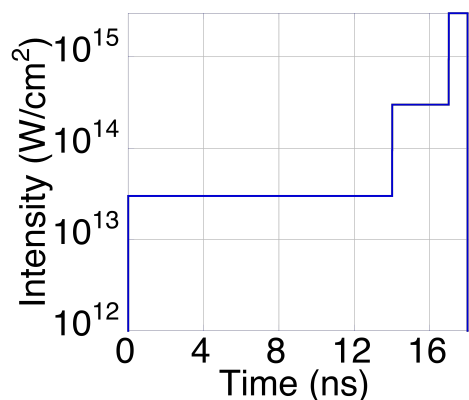
Laser Requirements:

Drive (8 quads): 3-5kJ/beam, focus= 1mm

XRTS Probe beams (12 quads): 75/beam (88ps impulse) or 4-5 kJ/beam (NIF Pulse), focus= 250 μ m

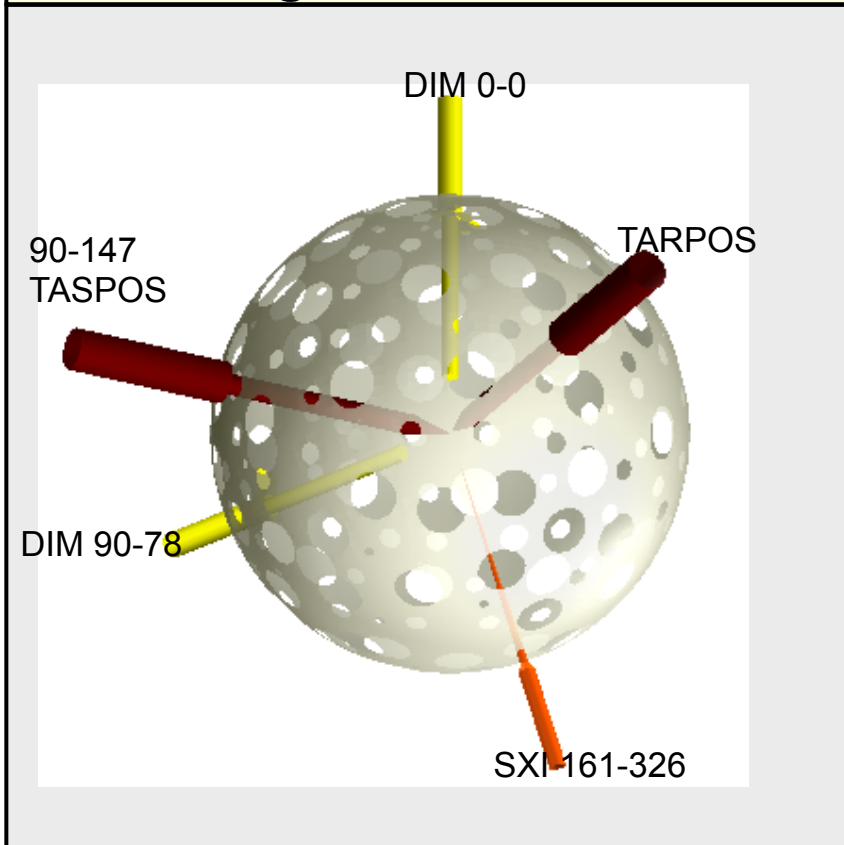
Radiography Backlighter (2 quads): 4-5 kJ/beam (NIF Pulse), focus= 1mm

Drive Pulse Shape: BL Pulse Shape (NIF):



Diagnostic configuration and compatible NIF platforms for planar high compression experiments

Experimental layout, w.r.t. target chamber



Diagnostics Configuration:

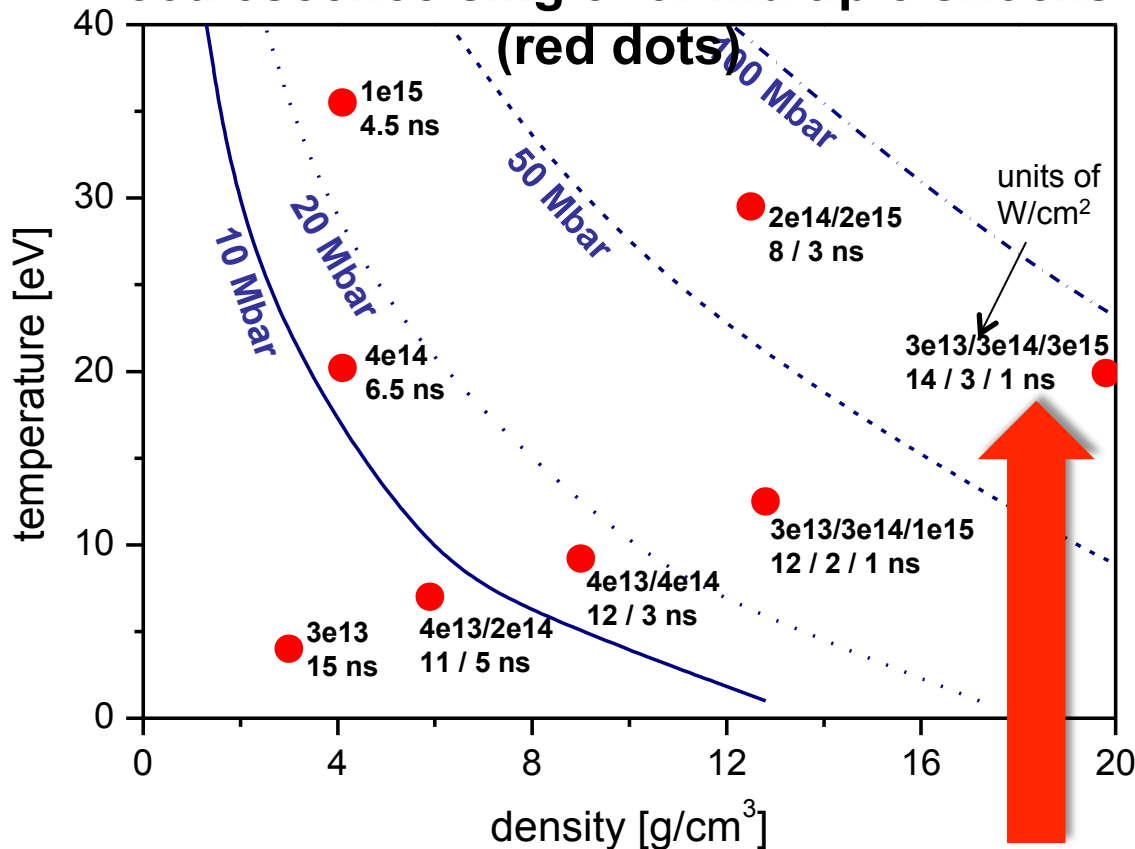
Diag	Location	Priority	Type	Calib
GXD or hGXI	90-78	Essential	3	Pre-Shot
hGXI or GXD +Supersnout 2 or MAHS	0-0	Essential	1	Pre-Shot
SXI 1	161-326	Ride-along	3	Pre-Shot

Compatible diag configurations:

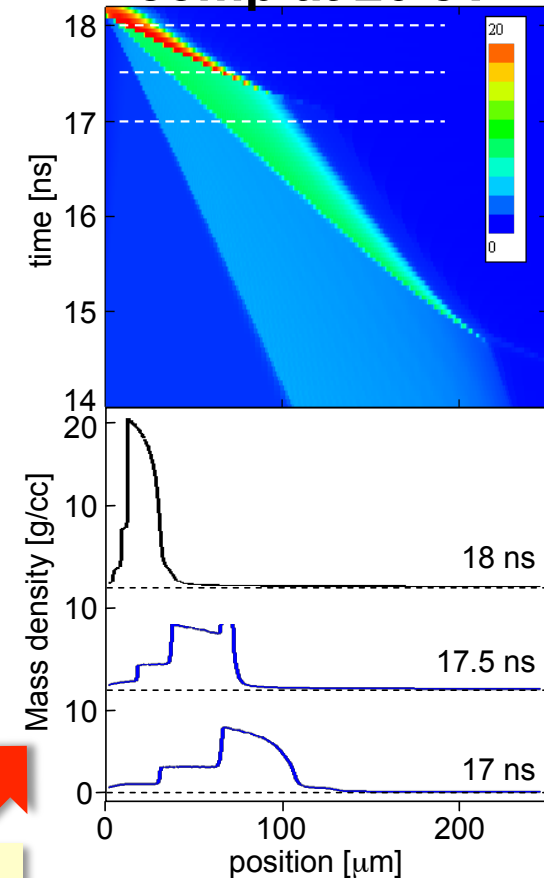
- 1) (DT4 for DT & ConAW) shown to the left
- 2) DT3 for DT & ConAW
- 3) DT2 for DT & keyhole

Pulse shaping enables creation of highly compressed and strongly coupled matter at relatively low temperatures

Densities and temperatures at shock coalescence single- or multiple shocks



We can reach 20x comp at 20 eV



3-shock comp

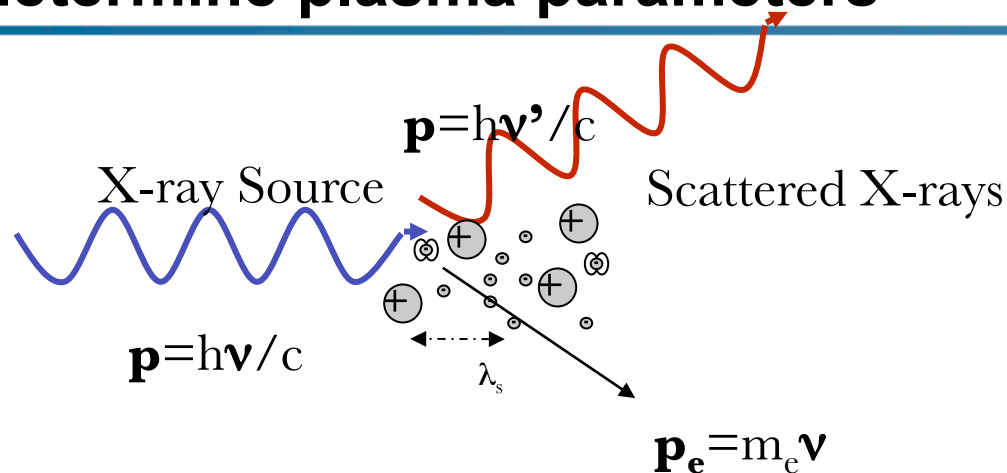
(PI) P. Neumayer, et al.

In these experiments X-rays are scattered from plasma electrons to determine plasma parameters

X-ray Scattering

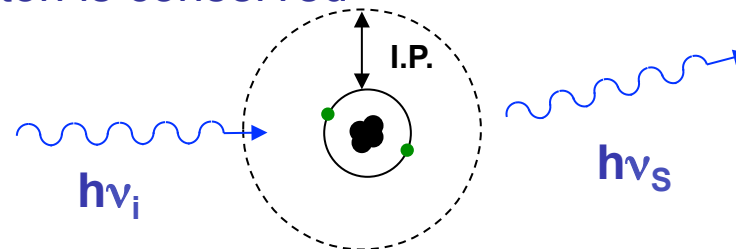
We Scatter x-rays from electrons in the plasmas.

Electrons absorb the photon, oscillate, and re-emit the x radiation.



Elastic (Rayleigh): E of incident photon is conserved

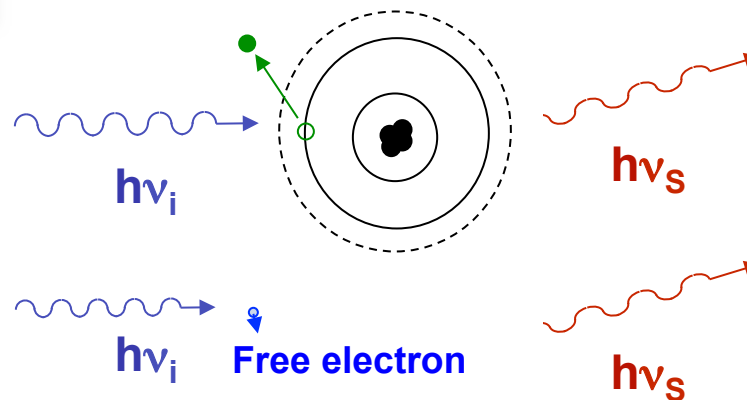
-Tightly Bound e^- :
Binding Energy > Compton Energy (ΔE_c)



Inelastic (Compton or Plasmon):

-Weakly Bound e^- :
Binding Energy < Compton Energy (ΔE_c)

-Free e^-



The non-collective and collective scattering will be applied to observe the micro and macroscopic motion of the electrons

Scattering Parameter: α

$$\alpha = \frac{1}{k\lambda_s} \propto \frac{\lambda}{\lambda_s}$$

$$k = \frac{4\pi}{\lambda_o} \sin(\theta/2)$$

Plasma Screening Length: λ_s

Fermi-degenerate Plasma:

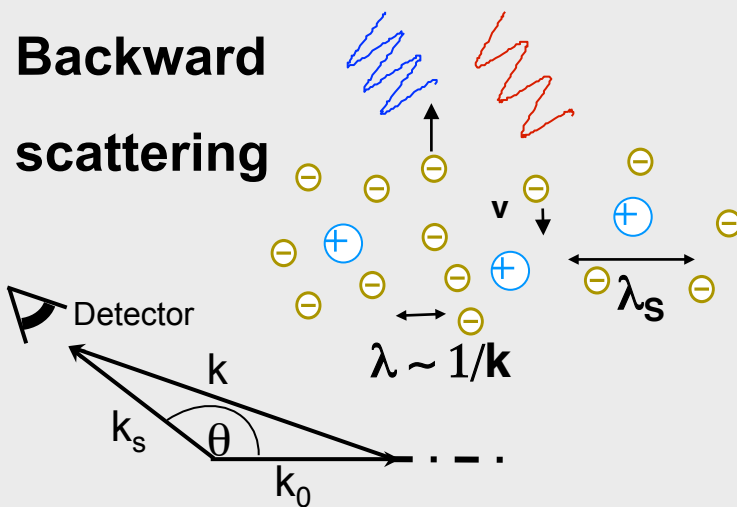
$$\lambda_{TF} = \left(\frac{\hbar^2}{4m_e e^2} \left(\frac{\pi}{3n_e} \right)^{1/3} \right)^{1/2}$$

Classical Plasma:

$$\lambda_D = \left(\frac{\epsilon_o k T_e}{n_e e^2} \right)^{1/2}$$

Non-Collective Scattering

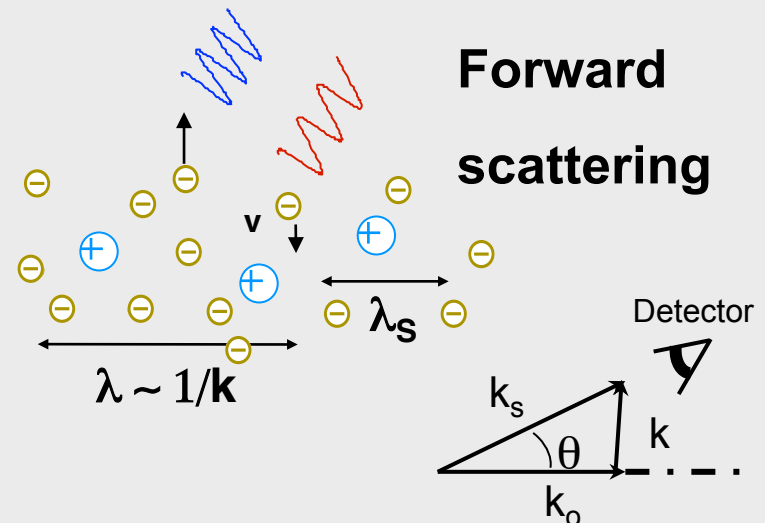
Probing: $\lambda < \lambda_s$ ($\alpha < 1$)



S. H. Glenzer et al., PRL (2003)

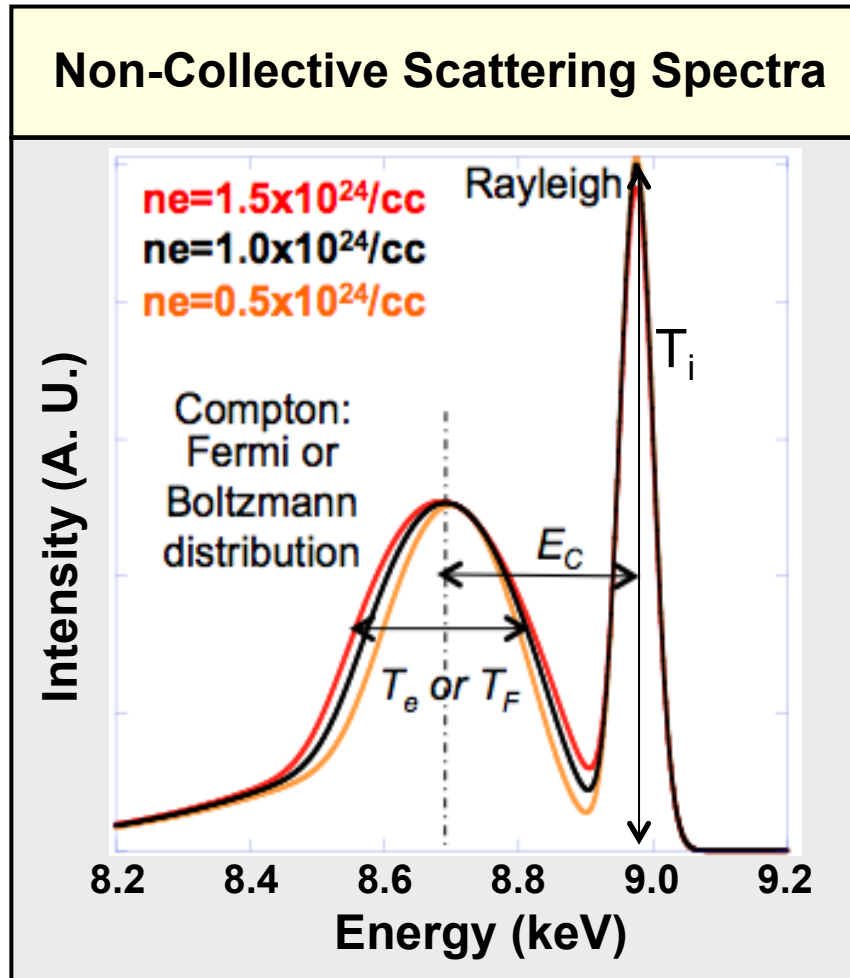
Collective Scattering

Probing: $\lambda > \lambda_s$ ($\alpha > 1$)

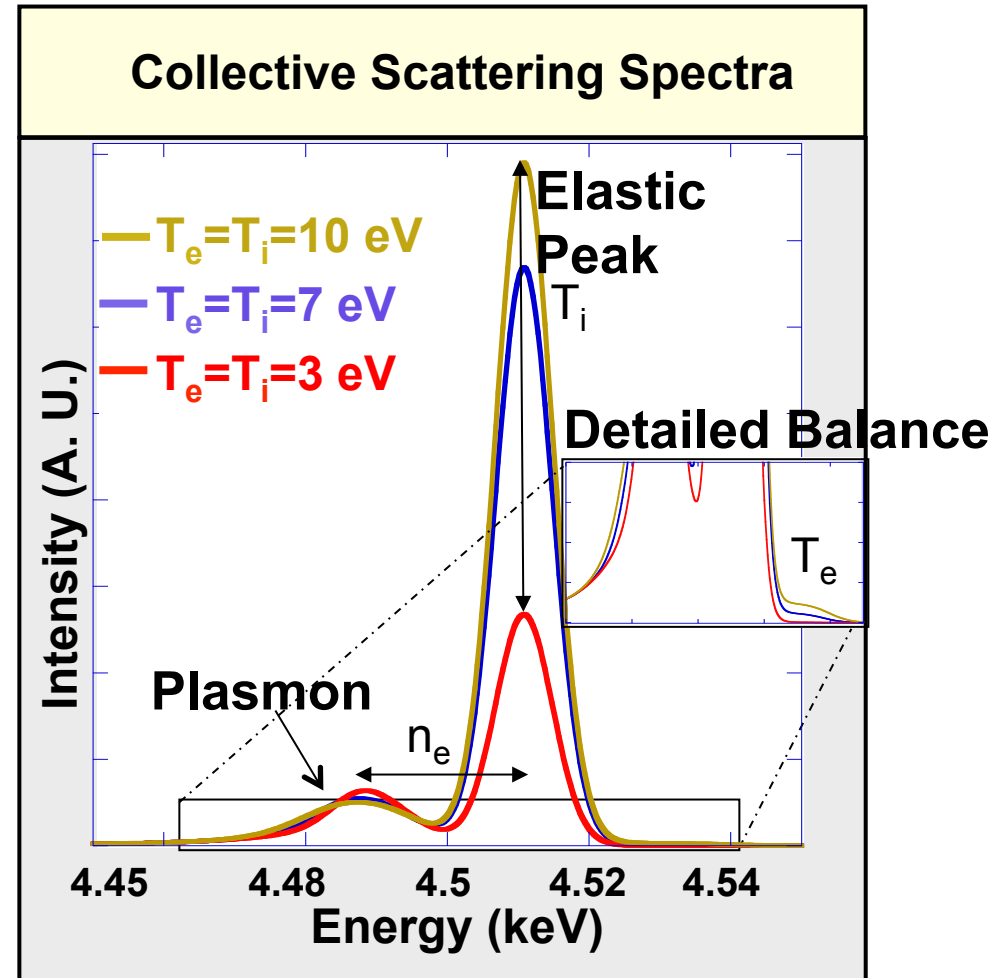


S. H. Glenzer et al., PRL (2007)

The plasma parameters can be determined from the shape of the scattered spectra



Partially Degenerate:
 Width $\rightarrow T_F (n_e)$
 Steepness of the red wing $\rightarrow T_e$
Non Degenerate:
 Width $\rightarrow T_e$



Detailed Balance: (Blue Feature) Take energy from plasmon wave $\sim e^{-\Delta\omega/kT}$

$$\Delta\omega \sim \omega_{pe} \propto n_e^{1/2}$$

The Intensity of the elastically scattered X-rays is Directly Related to the Structure Factors

Total Cross-Section Includes Free, Tightly, and Weakly bound States

$$\frac{d^2\sigma}{d\Omega d\omega} = \sigma_T \frac{k_1}{k_o} S(k, \omega)$$

$$S(k, \omega) =$$

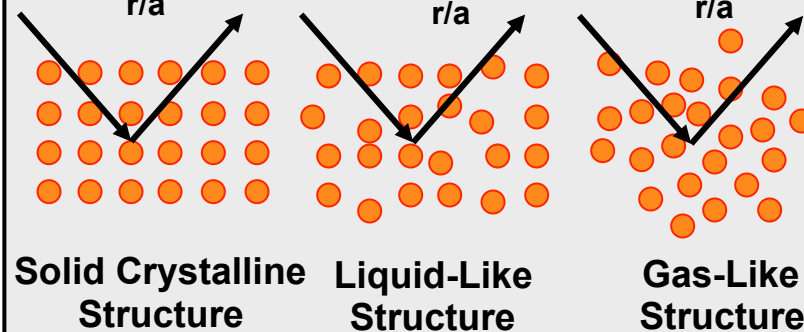
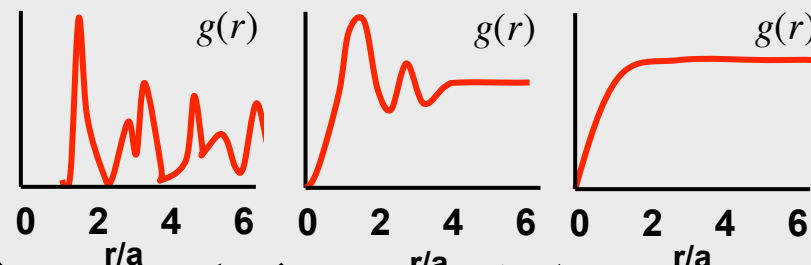
$$\underbrace{Z_f S_{ee}^o(k, \omega)}_{\text{Electron Feature}}$$

$$+ \underbrace{Z_c \int \tilde{S}_{ce}(k, \omega - \omega') S_{ce}(k, \omega') d\omega'}_{\text{Bound-Free Feature}}$$

$$+ \underbrace{|f_1(k) + q(k)|^2 S_{ii}(k, \omega)}_{\text{Ion Feature}}$$

The Scattering Intensity Depends on the Material Structure

$S_{ii}(k, \omega)$: Probability of finding an ion at a given distance from another ion (k space).

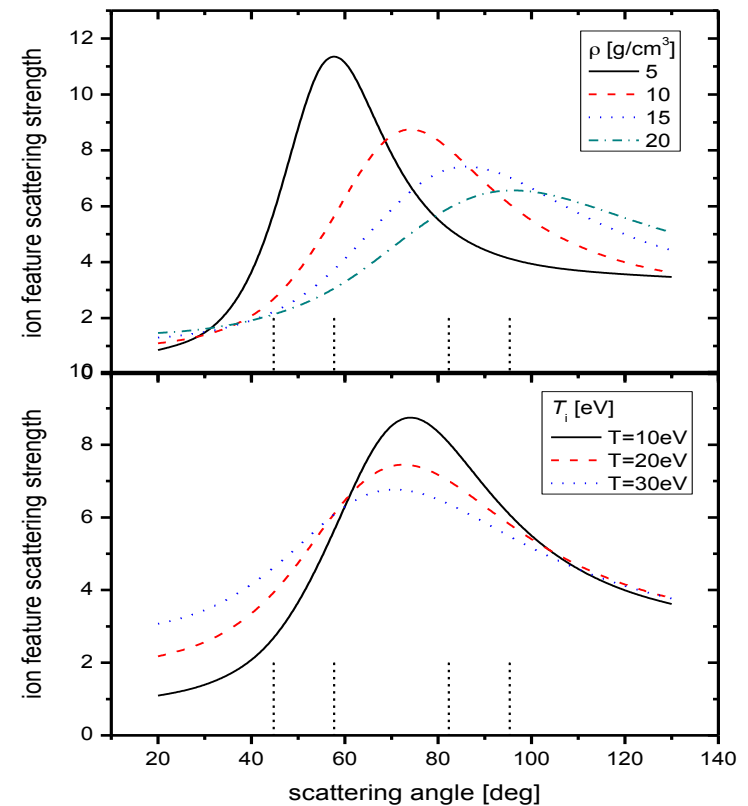


When atoms are less structured there is more of a probability to scatter at an arbitrary angle (less coupling) ==> $S(k, \omega)$ becomes smooth

By probing multiple scattering angles we can study ion structure in highly degenerate plasmas

- Study ion correlations via elastically scattered photons off of bound electrons (measure scattered intensity)
- Position of the correlation peak:
→ Wigner-Seitz radius
- Sharpness of the correlation peak:
→ Degree of ion-ion coupling
- Scattering at large k (17.5 keV) where $S_{ii} \sim 1$ will enable characterization of $|f(k) + q(k)|^2$

We can test models of material structure at high Temp & densities that predict ion-ion correlations and calculate EOS (DFT-MD)



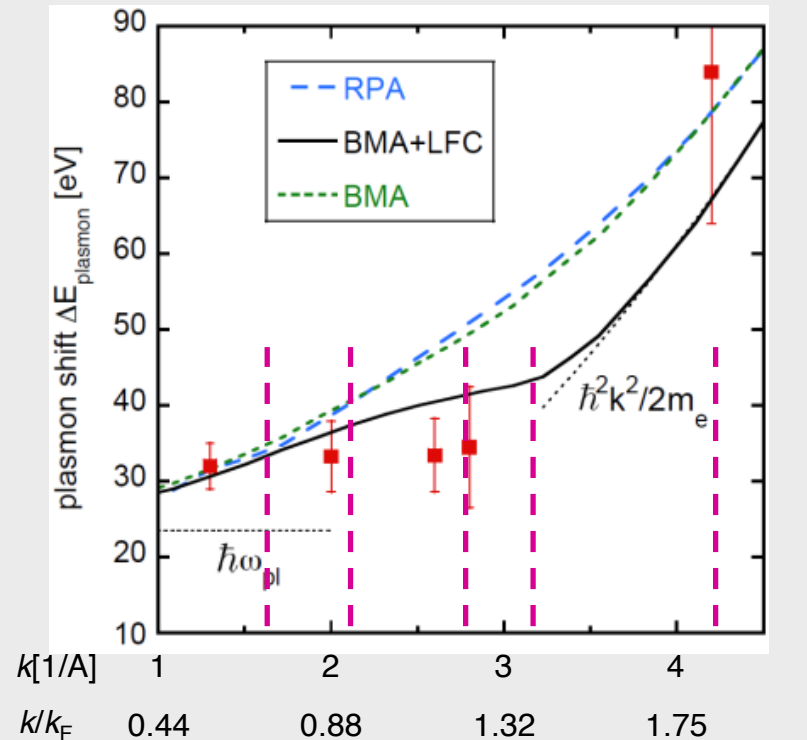
Neumayer (PI), Kritcher, Glenzer, Lee, et al

(PI) P. Neumayer, et al.

By probing multiple scattering angles we can study electron electron correlation in highly degenerate plasmas

- Collective scattering includes scatter of electron plasma waves (plasmons)
- RPA is used to calculate the plasmon dispersion
- Deviation from mean field theory (RPA) due to short range local field corrections for the electrons results in a reduced plasmon shift or kinetic energy (*e⁻ more correlated*)
- Dispersion measurements in dense plasmas are directly related to e⁻ coupling, i.e. internal energy

Energy shift of plasmon from the incident probe energy



— — — k/k_F at NIF, 20x compression, $Z_{\text{free}}=4$, 8.6keV backlighter

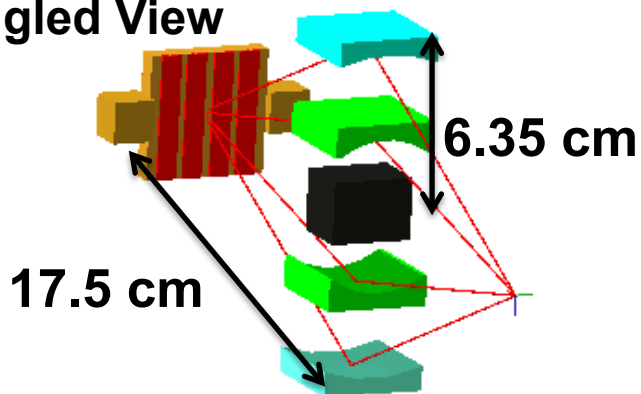
P. Neumayer et al., PRL 105, 075003 (2010)

(PI) P. Neumayer, et al.

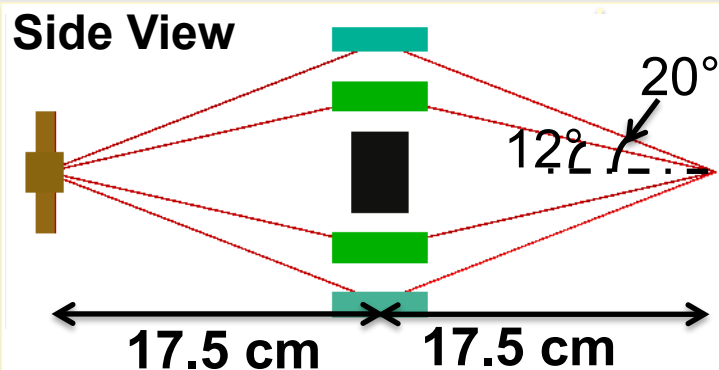
A new multi-angle spectrometer will enable simultaneous probing of four scattering angles or k-vectors

New Multi-angle XRTS Spectrometer

Angled View





Side View



Specs for spectrometer snout

- Angles: $\pm 20 \text{ deg} \pm 12 \text{ deg}$
- HOPG used for high reflectivity (3mrad), LiF is used for better energy and spatial resolution
- Crystals (7x5cm) will be cylindrically/ conically bent with $\Omega=7 \times 10^{-5}/\text{xtal}$
= 3×10^{-4} total
- Gated MCP's, IP surrounding MCP and full IP option
- Blast shield, filter options before the crystals and at the MCP, Tungsten block for direct line of sight of TCC to the MCP

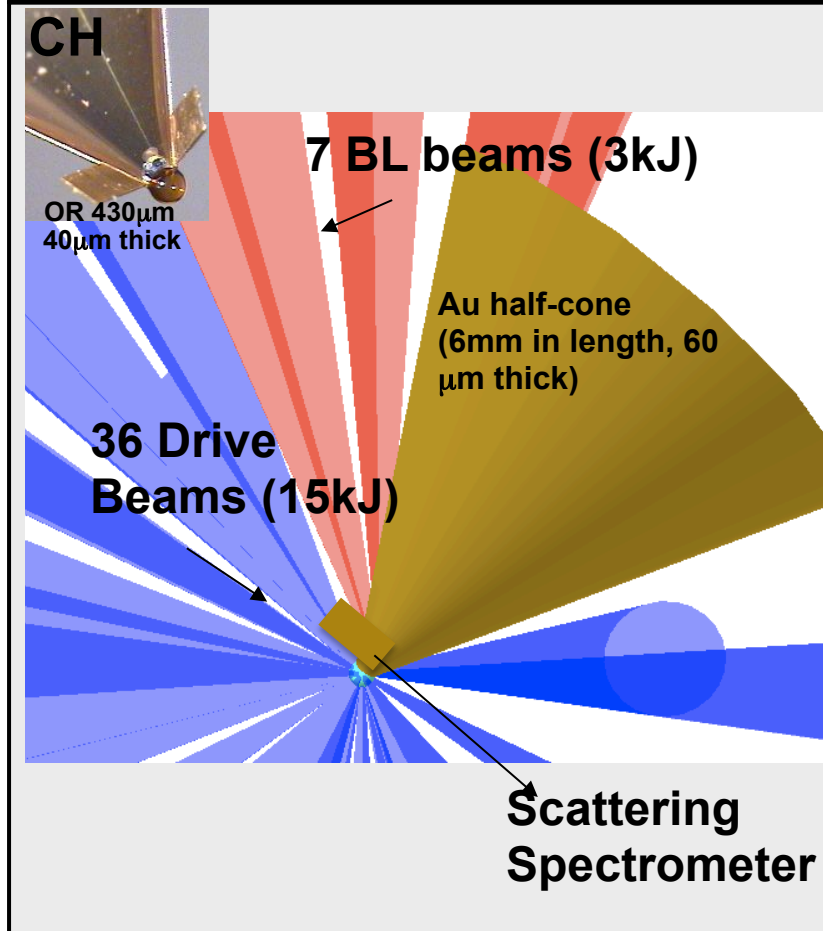
 LiF(200) ROC: 60 mm, 7.4-10.6 keV, first order

 HOPG ROC: 38 mm, 7.3-10.7 keV, first order

Fabrication through Artep.

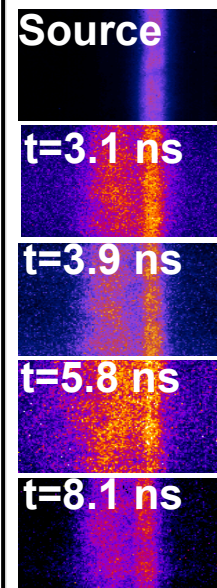
Highly compressed and degenerate matter has been probed with XRTS from spherically imploded capsules at Omega

Schematic of the experimental configuration

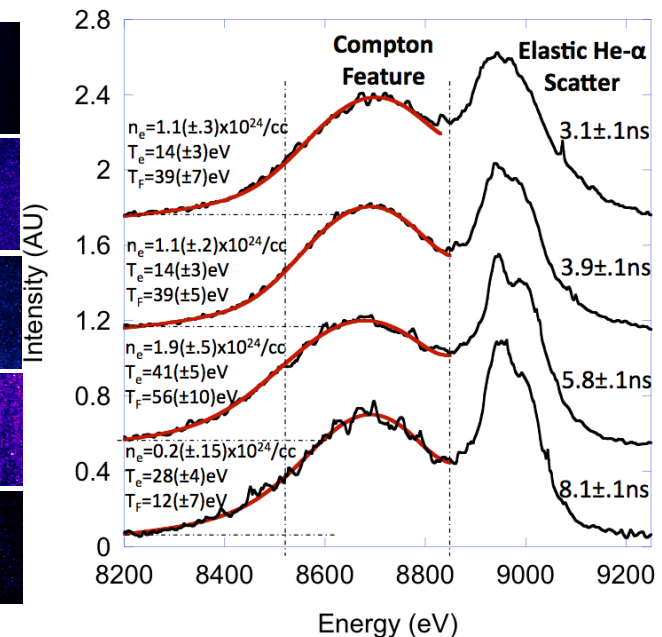


Compressions of x7 and x13 times solid density in CH and Be have been probed

Raw data



Profiles of the experimental data



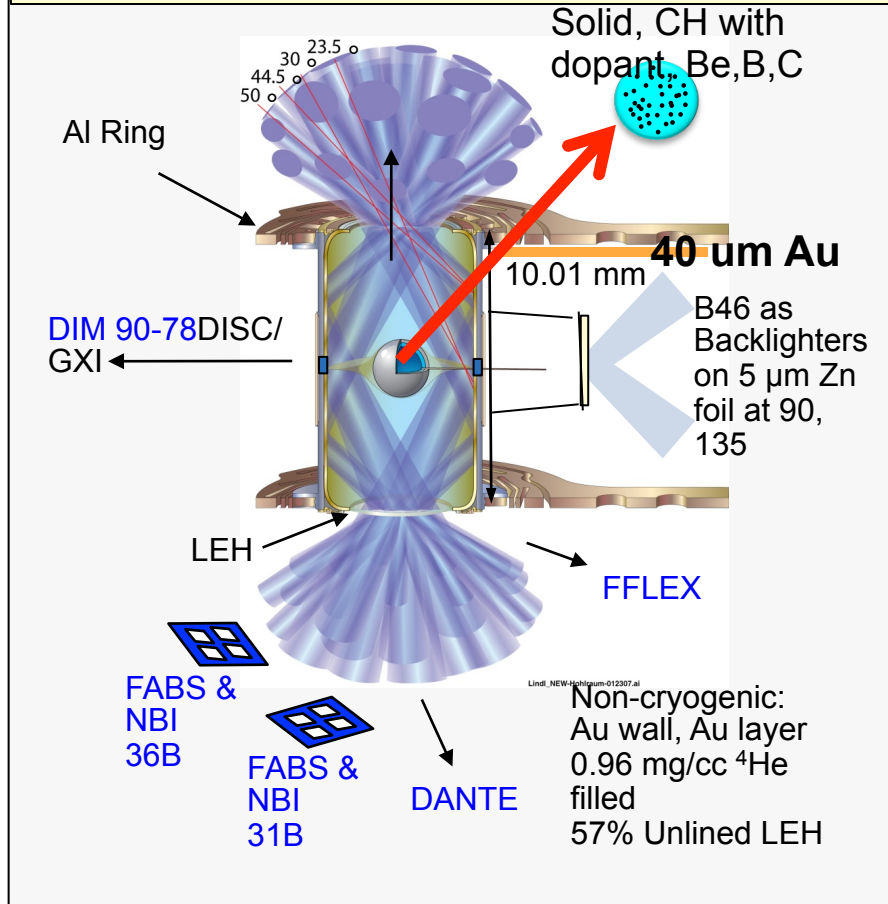
SNR of up to >200 (A. krichter, et al. PRL)

→ Infer plasma properties such as adiabat and degree of ion-ion coupling (>5)

The high SNR of the data and sensitivity of the Compton-red wing enabled model independent determination of T_e and n_e (~20% for n_e and T_e , and 13% for adiabat)

NIF enables the creation of conditions to study the cores of exoplanets (up to 2Gbar) and brown dwarfs using a fielded NIC platform

Solid spheres will be indirectly driven in a spherically convergent geometry



Closely related to the NIC Con-A platform
(D. Hicks)

(PI) R. Falcone, et al.

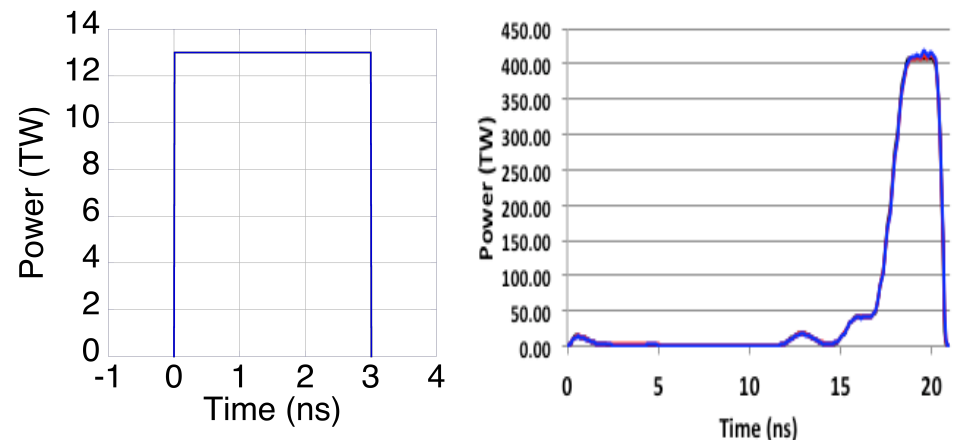
Laser Requirements:

Drive : 6.9kJ/beam, 1.3MJ total (NIF ignition pulse (575)), focus= 1mm

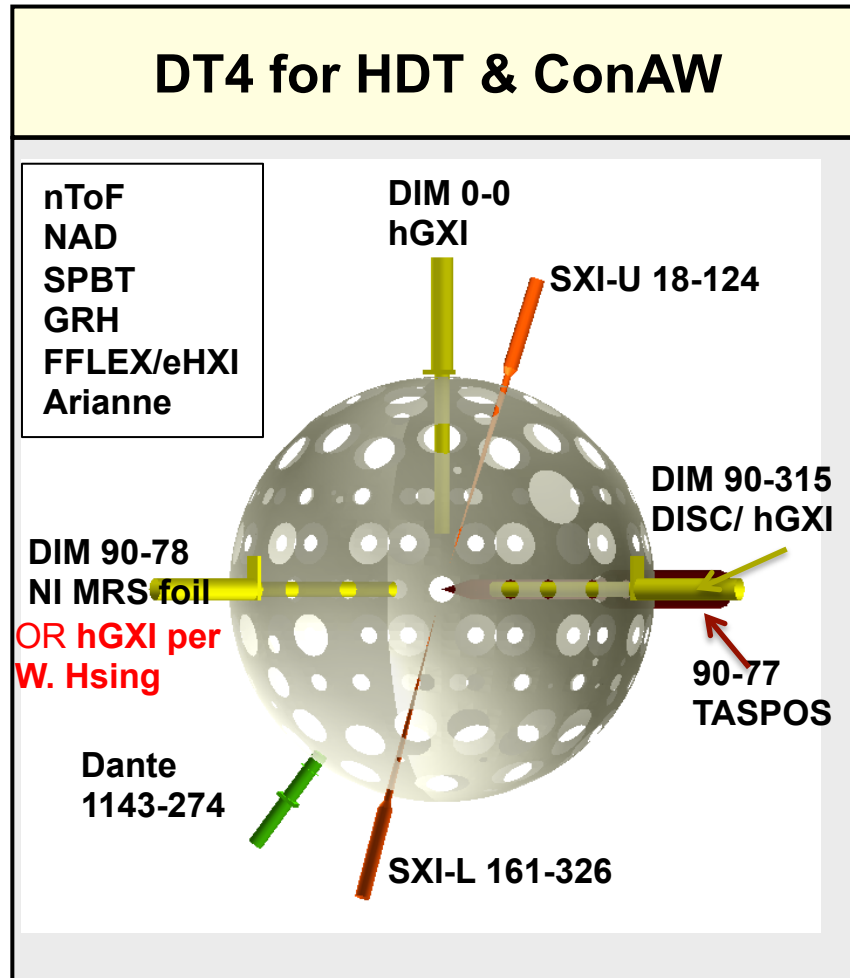
XRTS probe and Radiography

Backlighter (2 quads): 4-5 kJ/beam (NIF ignition pulse (575)), focus=500µm

BL (3ns) and Drive Pulse Shape (NIF):



Diagnostic configuration and compatible NIF platforms for Gbar pressure experiments



Diagnosics Configuration:

Diagnostic list	Location	Priority	Type	Calib
GXD or hGXI	90-315 or 90-78	Essential	1	Pre-Shot
hGXI +Supersnout 2 or MAHS	0-0	Essential	1	Pre-Shot
SXI 2	18,123	Ride-along	3	Pre-Shot
SXI 1	161,326	Ride-along	3	Pre-Shot
Dante 1	143,274	Ride-along	3	Pre-Shot

Other compatible diag configurations:

- 1) (DT4 for DT & ConAW) shown to the left
- 2) DT3 for DT & ConAW
- 3) DT2 for DT & keyhole

The experimental choice of solid sphere target depends on transmission through the target and changes in opacity

Compromise:

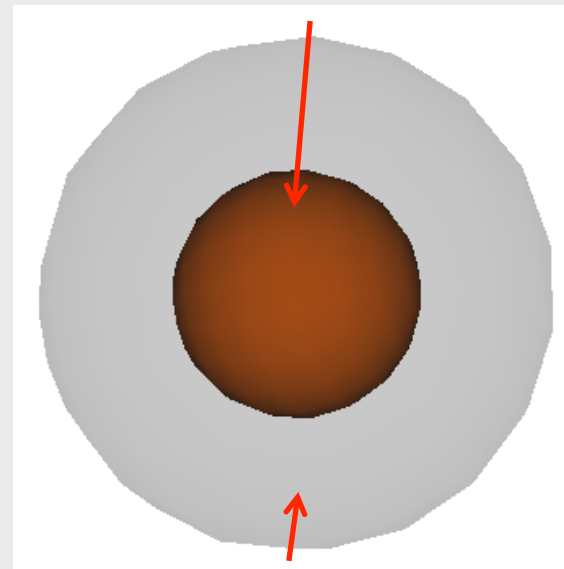
- Transmission through sample (small) vs spatial resolution (large).
- Compressible: more heating => worry about opacity falling.

Best: C(diamond) with thick “ablator”

Possible: C(graphite), BN, B₄C, with ablator; B or Be with or without ablator, doped CH

Solid spheres are used instead of NIC capsules to better constrain the mass density through profile fitting

Sample: diamond or graphite, 1 mm diameter



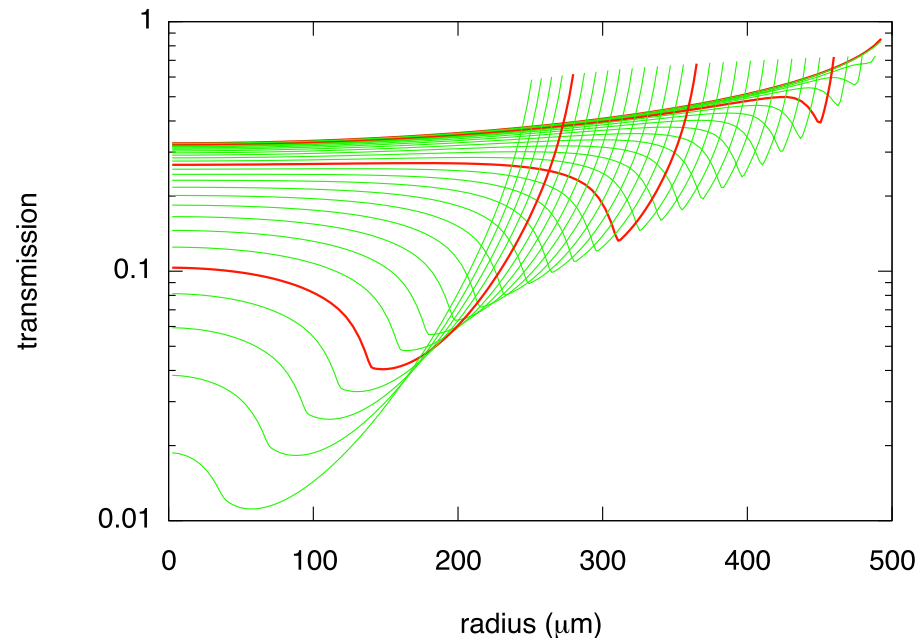
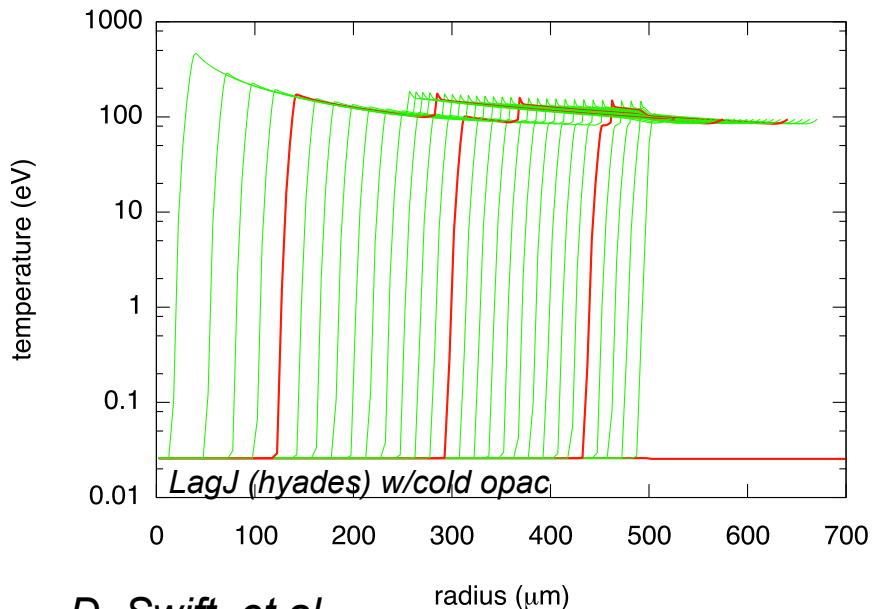
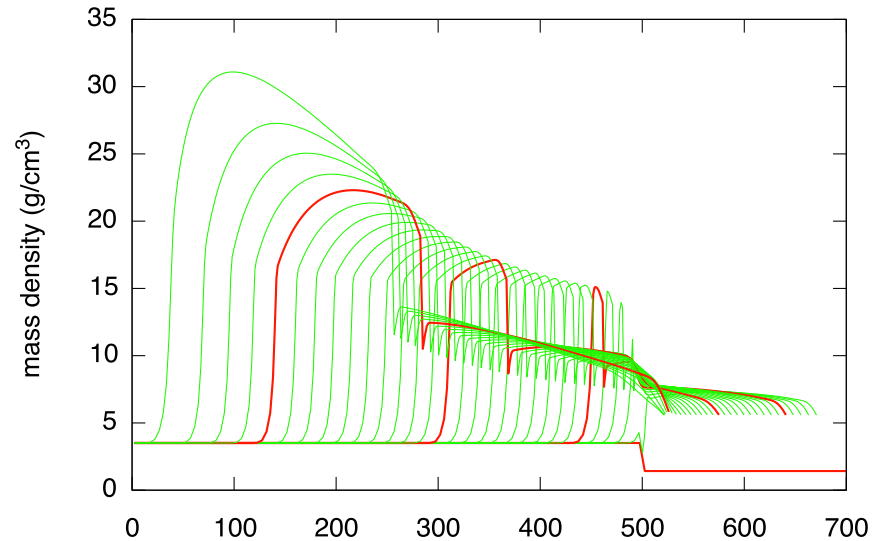
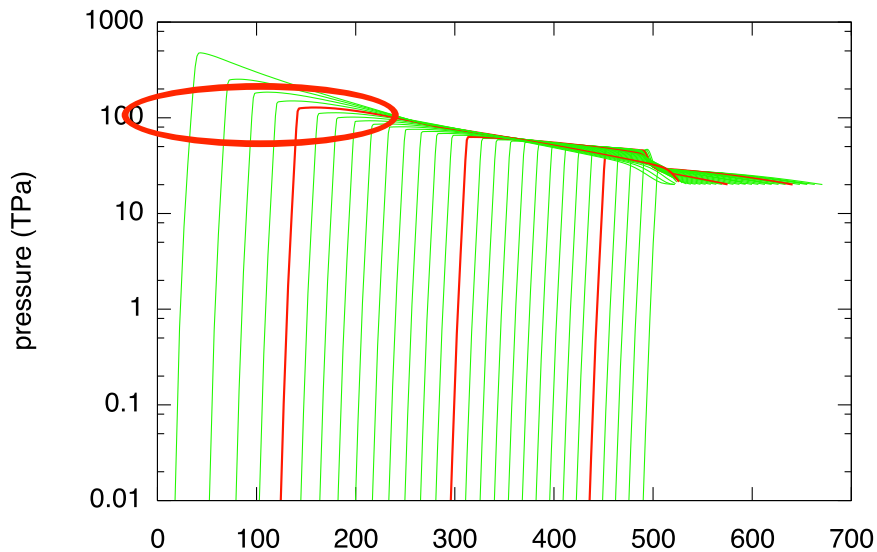
Ablator: CH, 2 mm diameter, doped with Ge/Br

D. Swift, et al.

Diamond is also important for studying the structure of giant planets and advanced ICF ablator materials.

Rad-hydro simulations show that we can reach >2 Gbar pressures for indirectly driven solid* targets

*Carbon (dia) with thick ablator, 20 TPa drive

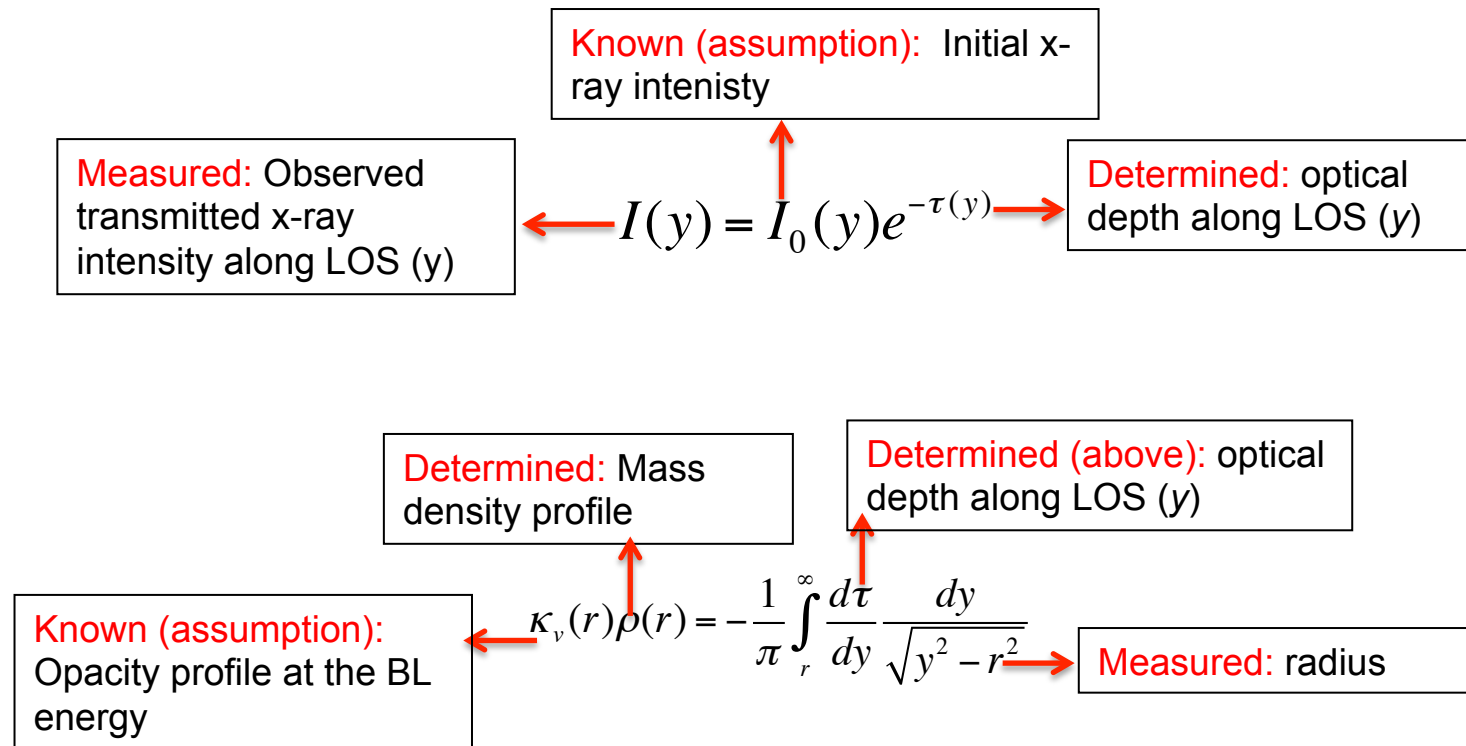


C EOS: SESAME #7830

D. Swift, et al.

A mass density profile of the sample can be determined from a radiographic image using Abel inversion

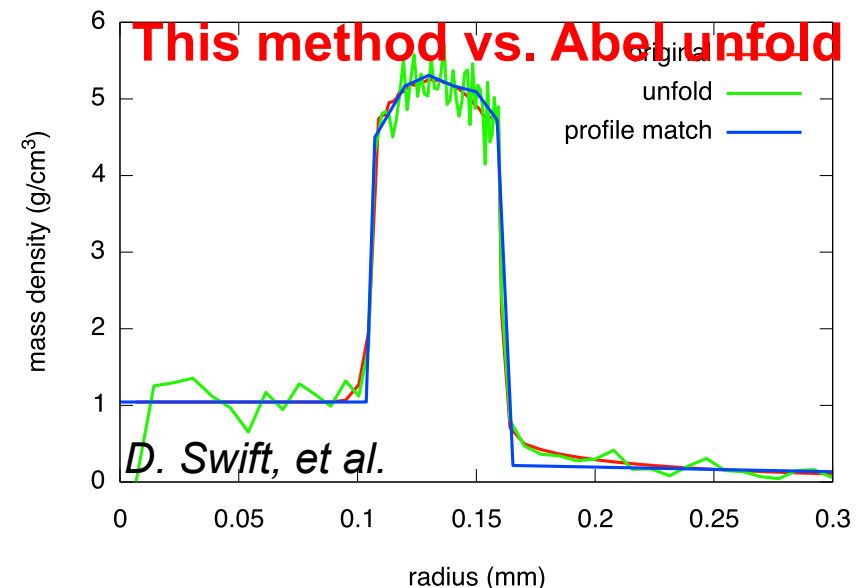
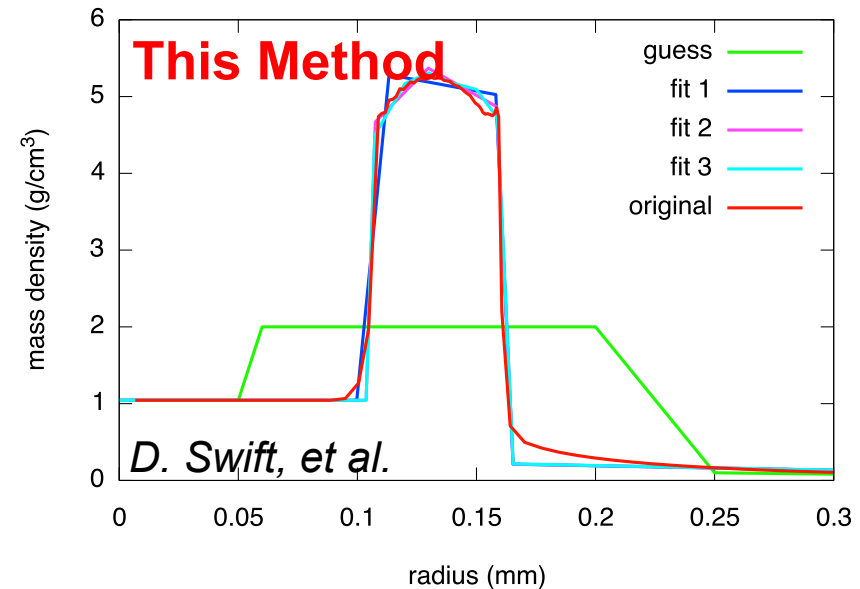
Abel inversion: An integral transform that can be used to reconstruct a mass density profile from a projected radiographic image.



However... in addition to the assumptions, Abel inversion can amplify noise from photon statistics and backgrounds that is already an issue

We plan to use solid samples and profile matching with Bayesian analysis to determination the mass density profile

- **Bayesian analysis:**
 - Use a density profile with adjustable parameters (e.g. nodes) and constraints
 - Simulate radiograph
 - Adjust parameters for best fit
 - Refine
- **Solid Target:** Using known unshocked density, captures shock jump much more accurately, even with noise and inaccurate profile of outer layers



With the mass density profile determined and measured shock speeds we can constrain the EOS

Rankine-Hugoniot Relations:

$$D^2 = v_0^2 \frac{p - p_0}{v_0 - v}$$

$$u^2 = (p - p_0)(v_0 - v)$$

$$e = e_0 + \frac{1}{2}(p + p_0)(v_0 - v)$$

D = Shock speed

$v = \frac{1}{\rho}$ = Specific volume

p = Pressure

u = Particle speed

e = Internal energy

Re-arrange D^2 relation:

$$p = p_0 + \frac{D^2}{v_0^2}(v_0 - v)$$

$$\frac{\delta p}{p} = \frac{\delta(v_0 - v)}{v_0 - v} + 2 \frac{\delta D}{D} \sim \frac{\delta \rho}{\rho} + 2 \frac{\delta D}{D}$$

Only depends on:

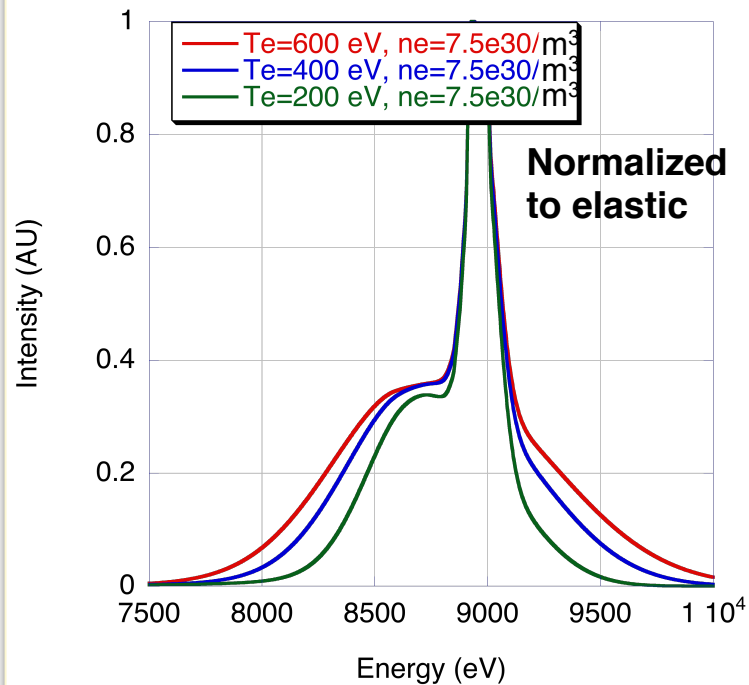
- D : Shock speed (measured)
- ρ : Mass density (determined)
- P : **Pressure** (inferred)

R. Falcone, D. Swift, et al.

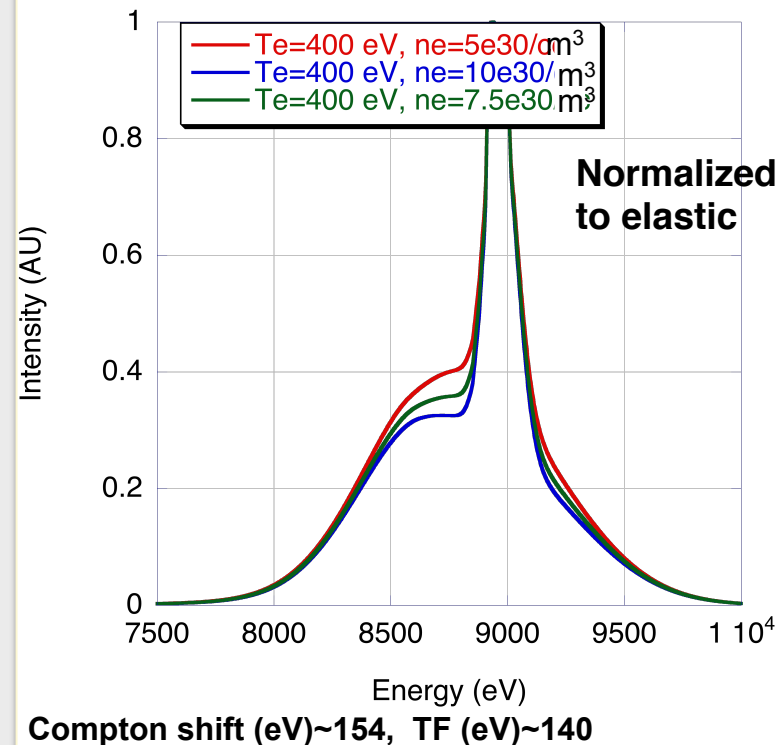
→ With P and ρ we determine the equation of state at Gbar pressures for CH and Diamond (*Other materials under consideration: Be, B, doped CH*)

For non-degenerate plasmas the width of the Compton scattered feature is sensitive to T_e

Sensitivity of x-ray scattering to electron temperature



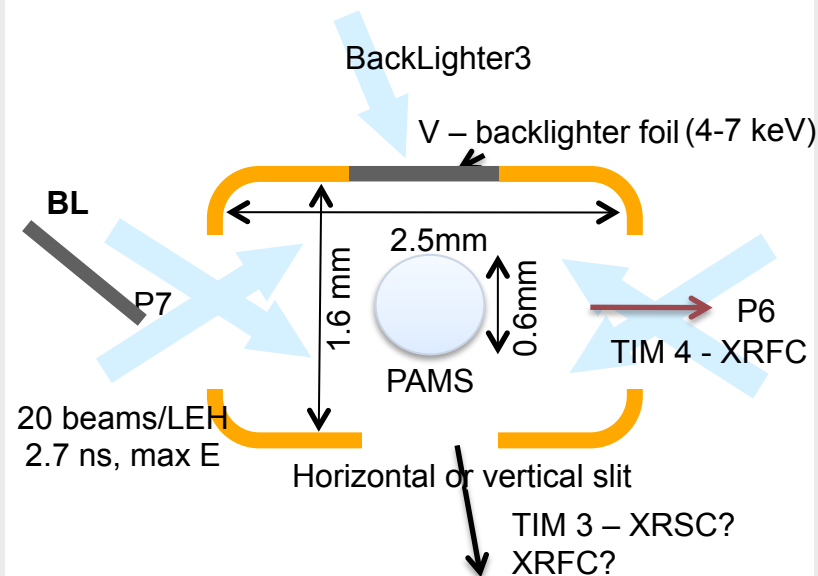
Sensitivity of x-ray scattering to changes in electron density



XRTS: is sensitive to temperature due to Doppler broadening in the non-collective regime at these temperatures and densities

Current radiography experiments at Omega using solid spheres observe the shock front and initial densities have been estimated

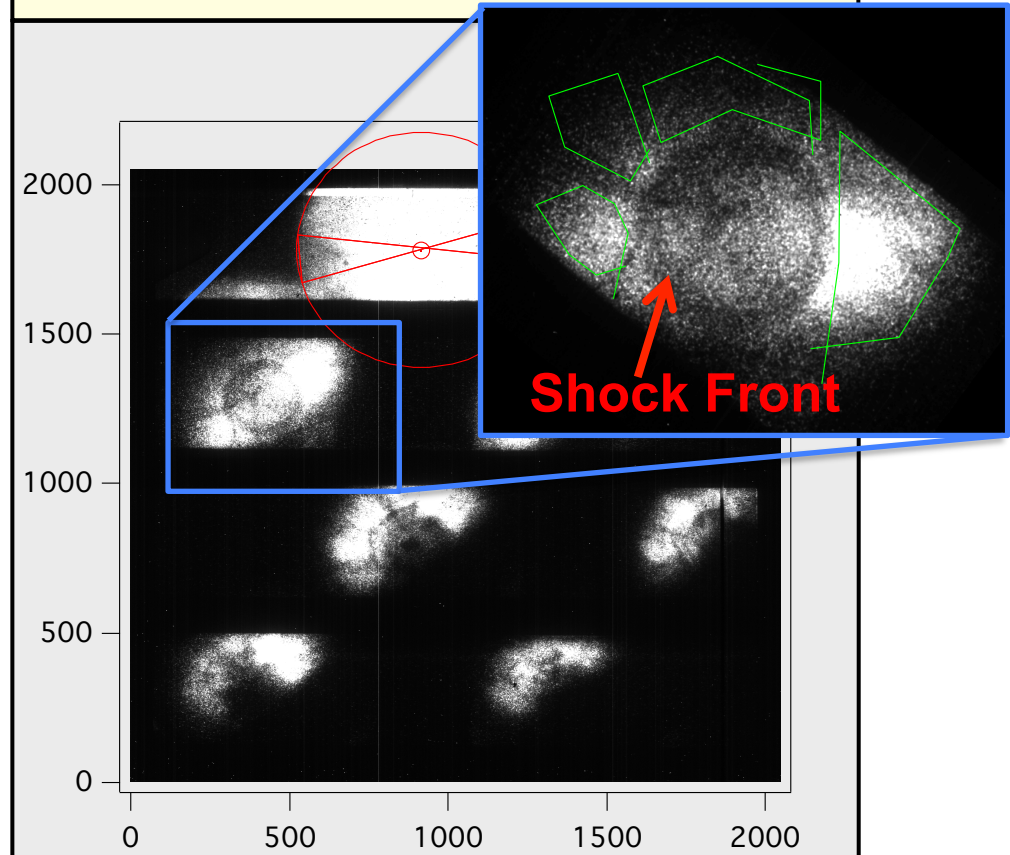
Supporting Omega experiments are currently being fielded



T4-XRFC (2D radiograph), DANTE (Tr), T3-XRS (BL), SSCA (streaked radiograph)

J. Hawreliak, et al.

Measured radiographs of indirectly driven solid PAMS sample at Omega



Preliminary: Meas. were in rough agreement with predicted shock velocities and shell thickness (LASNEX); inferred densities were lower than predicted

Summary and outlook

- We have shown that the n_e and T_e for plasmas of densities of up to x13 times solid density can be accurately characterized with XRTS
- At NIF we plan to probe highly degenerate and compressed plasmas (20eV and 20x compression) and plasmas at high pressure and temperature (>1Gbar and ~0.4 keV)
- Other supporting Omega experiments probe the EOS of CH in a solid-sphere geometry (*J. Hawreliak*) and develop high Z backlighters such as Mo K-a and He-a (*T. Ma*)
- There have been many more XRTS and radiography experiments at Omega and Titan that investigate ion structure in 2x or 3x solid density plasmas at low temperatures

The end.



NIF

